

# Witnessing Entanglement with the Stabilizer Formalism

Géza Tóth (MPQ)

Otfried Gühne (Innsbruck)

**quant-ph/0501020**

**PRL 94, 060501 (2005)**

DPG, Berlin, March 2005

# Outline

- Posing the problem.
- Stabilizing operators for GHZ and cluster states
- Entanglement conditions with stabilizing operators
- Characterizing entanglement conditions: Noise tolerance
- Criteria for detecting genuine multi-qubit entanglement (or estimating the fidelity)
- **Nonlinear entanglement criteria**

# Main problems in a *many*-qubit system

- Let us consider systems in which the qubits are individually accessible, but only *local* measurements are possible.
- Complete state tomography is very hard since the number of measurements increases *exponentially* with the number of qubits.
- Measuring a usual entanglement witness or the fidelity with respect to a given state is also very hard.
- We look for a solution using stabilizer theory. This makes it possible to characterize the state with few measurements.

# Stabilizing operators

- An operator  $S_k$  *stabilizes* the quantum state  $|\Psi\rangle$  if

$$|\Psi\rangle = S_k |\Psi\rangle$$

- Many quantum states can be more efficiently described by their stabilizing operators, than by the state vector.
- This is used in error correction [Gottesman PRA 96].

# Stabilizing operators for the GHZ state

- Example: GHZ state:  $|000\rangle + |111\rangle$

$$S_1 = X^{(1)} X^{(2)} X^{(3)},$$

Stabilized by:  $S_2 = Z^{(1)} Z^{(2)},$

$$S_3 = Z^{(2)} Z^{(3)}.$$

- The GHZ state is *uniquely* characterized by

$$\langle S_1 \rangle = \langle S_2 \rangle = \langle S_3 \rangle = +1.$$

# Our first sufficient entanglement condition

- For pure product states

$$\begin{aligned}\langle S_1 + S_2 \rangle &= \langle X^{(1)} X^{(2)} X^{(3)} + Z^{(1)} Z^{(2)} \rangle \\ &= \langle X^{(1)} \rangle \langle X^{(2)} \rangle \langle X^{(3)} \rangle + \langle Z^{(1)} \rangle \langle Z^{(2)} \rangle \\ &\leq \left| \langle X^{(1)} \rangle \langle X^{(2)} \rangle \right| + \left| \langle Z^{(1)} \rangle \langle Z^{(2)} \rangle \right| \leq 1\end{aligned}$$

- Due to the convexity of separable states, this is also true for mixed separable states.

# Our first sufficient entanglement condition II

- For separable states

$$\left\langle X^{(1)} X^{(2)} X^{(3)} + Z^{(1)} Z^{(2)} \right\rangle \leq 1$$

If this bound is violated then the state is entangled.

- For the GHZ state

$$\left\langle X^{(1)} X^{(2)} X^{(3)} + Z^{(1)} Z^{(2)} \right\rangle = 2$$

# Stabilizing operators for cluster states

- Stabilizing operators for an  $N$ -qubit cluster state

$$S_1 = X^{(1)} Z^{(2)},$$

$$S_k = Z^{(k-1)} X^{(k)} Z^{(k+1)}; \quad k = 2, 3, \dots, N-1,$$

$$S_N = Z^{(N-1)} X^{(N)}.$$

- For separable states

$$\langle S_k + S_{k+1} \rangle \leq 1$$

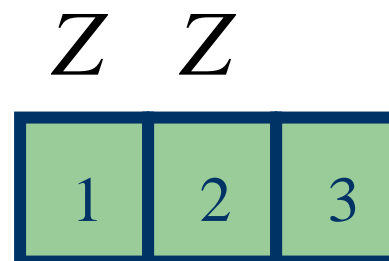
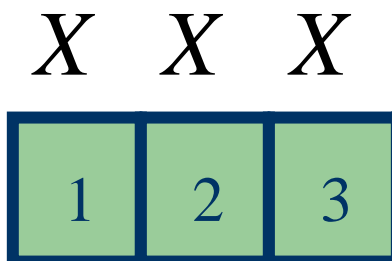


## How to measure the conditions?

- Advantage of our conditions that they are **easy to measure locally**

$$\left\langle X^{(1)} X^{(2)} X^{(3)} + Z^{(1)} Z^{(2)} \right\rangle \leq 1$$

- *Two local measurement settings are needed*



## Noise tolerance

- In an experiment the GHZ state is never prepared perfectly

$$\rho = (1 - p_{noise}) |GHZ_3\rangle\langle GHZ_3| + p_{noise} \rho_{completely\_mixed}$$

- For each entanglement condition there is a noise limit. For a noise larger than this limit the GHZ state is not detected as entangled.

## Noise tolerance II

- Let us take the condition

$$\langle S_1 + S_2 \rangle = \langle X^{(1)} X^{(2)} X^{(3)} + Z^{(1)} Z^{(2)} \rangle \leq 1$$

Noise is tolerated if  $p_{noise} < 1/2$

- Better condition, tolerating noise if  $p_{noise} < 2/3$

$$\langle S_1 + S_2 + S_1 S_2 \rangle =$$

$$\langle X^{(1)} X^{(2)} X^{(3)} + Z^{(1)} Z^{(2)} - Y^{(1)} Y^{(2)} X^{(3)} \rangle \leq 1$$

# Genuine multi-qubit entanglement

- Genuine three-qubit entanglement

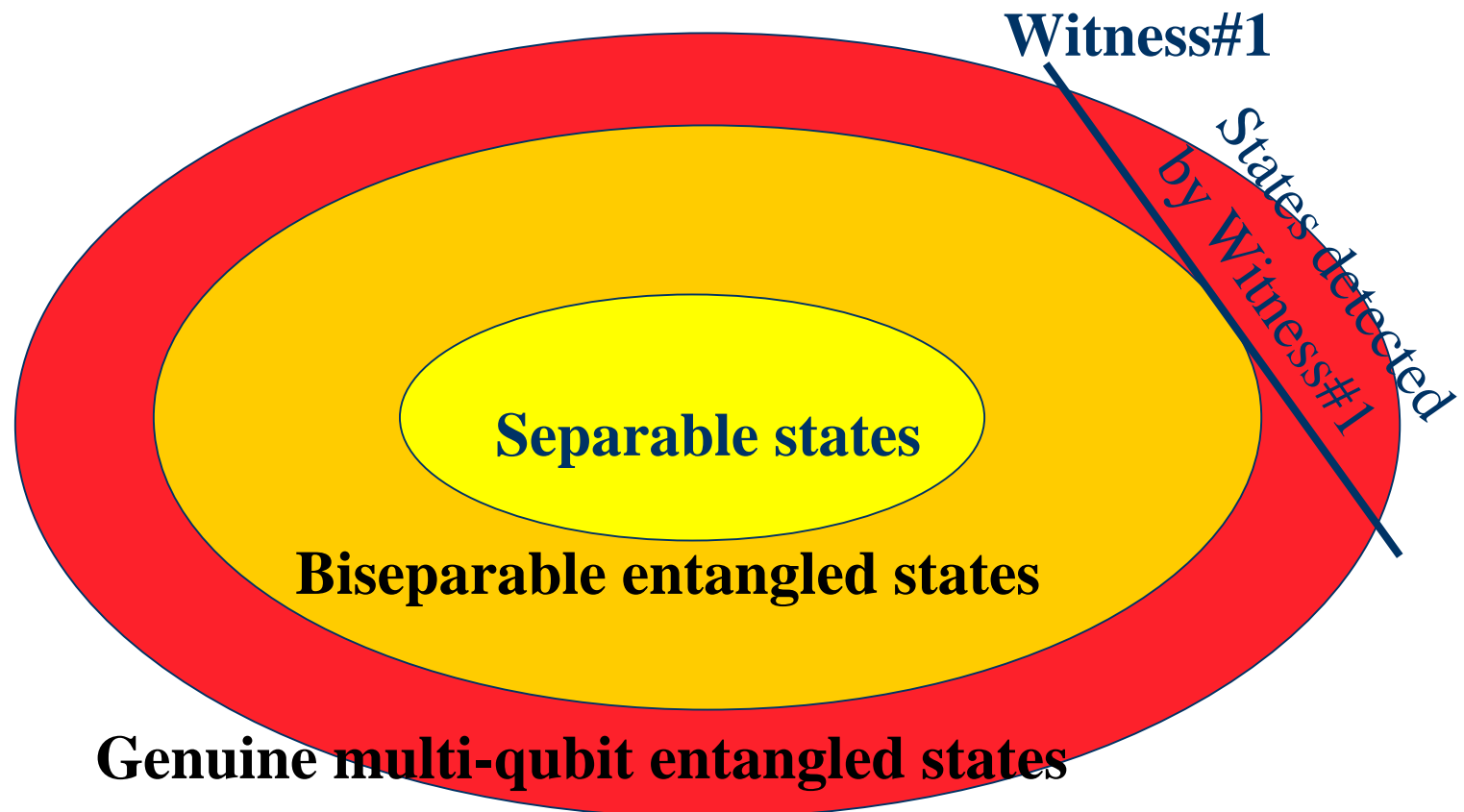
$$|000\rangle + |111\rangle$$

- Biseparable entanglement

$$|001\rangle + |111\rangle = (|00\rangle + |11\rangle)|1\rangle$$

- A mixed entangled state is *biseparable* if it is the mixture of biseparable states (of possibly different partitions).

# Entanglement witnesses for detecting genuine multi-qubit entanglement



## Witnesses for detecting genuine multi-qubit entanglement

$$W_{GHZ_N} = 3 - 2 \left[ \frac{1 + S_1^{(GHZ_N)}}{2} + \prod_{k>1} \frac{1 + S_k^{(GHZ_N)}}{2} \right]$$



- Advantage: only the minimal two measurement settings are needed, no exponential increase with the number of qubits!
- Similar ideas can be used to estimate the fidelity.

# Conditions based on uncertainty relations

- Based on the ideas of Gühne PRL 92, 117903 for biseparable states of the type (1)(23) we have

$$\left\langle X^{(1)} X^{(2)} X^{(3)} + Z^{(1)} Z^{(2)} \right\rangle + \frac{1}{2} \left( \left\langle X^{(1)} + X^{(2)} X^{(3)} \right\rangle^2 + \left\langle Z^{(1)} + Z^{(2)} \right\rangle^2 \right) \leq 1$$

- Nonlinear terms: play the role of “refinement“

## Conditions based on entropic uncertainty relations

- Based on the ideas of [Gühne & Lewenstein PRA 70, 023316] for biseparable states we have,

$$\sum_{k=1}^N H(S_k) \leq \ln 2$$

where  $H(A)$  is the Shannon entropy of the outcomes after measuring operator  $A$ .



**PRL 94, 060501 (2005)**

**quant-ph/0501020**

## Summary

- We constructed entanglement conditions with stabilizing operators.
- Our conditions detect entangled states close to GHZ and cluster states.
- Home page:  
<http://www.mpq.mpg.de/Theorygroup/CIRAC/people/toth>
- \*\*\*\*\* THANK YOU!!! \*\*\*\*\*